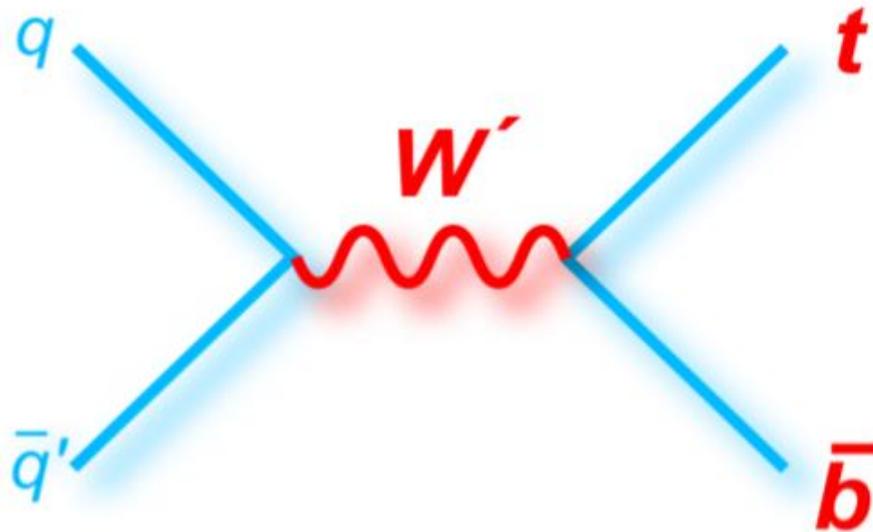


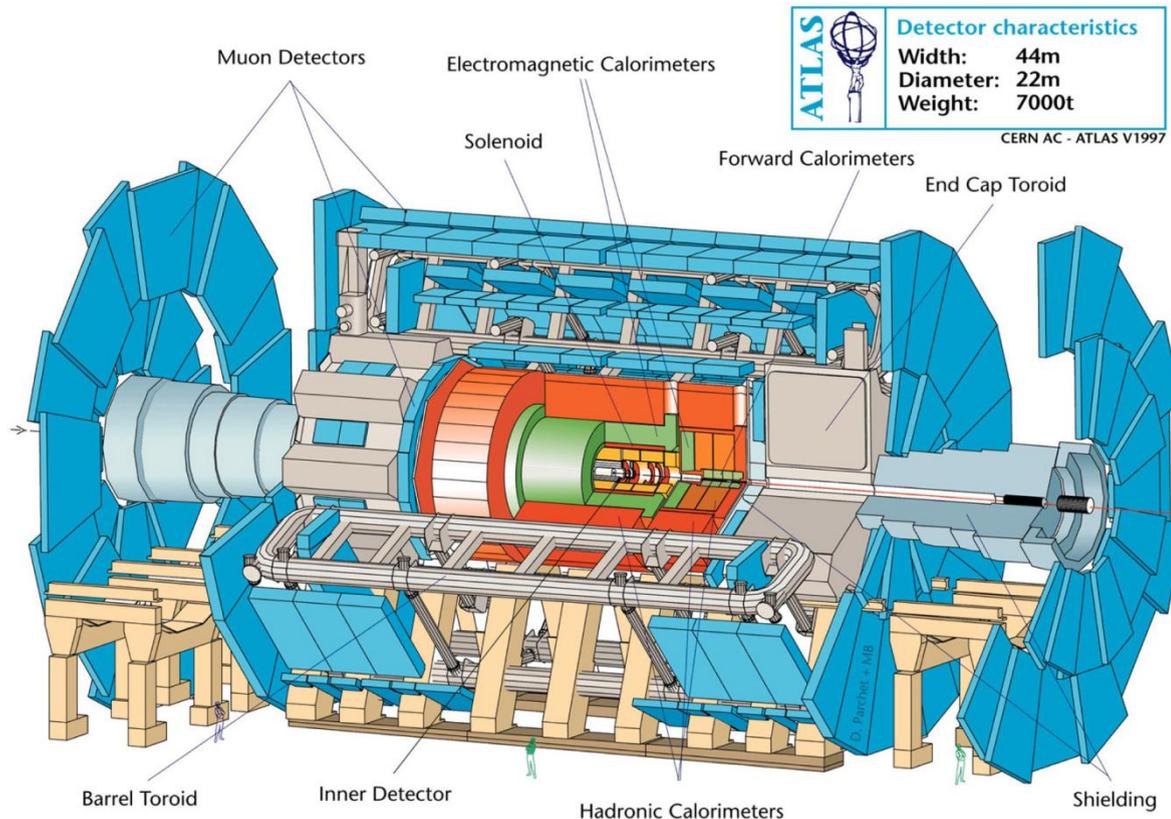
# Search for $W'$ Production in the Single Top Channel with the ATLAS Detector



**Patrick True (MSU)**  
**August 15<sup>th</sup> 2013**

- I am presenting recent results from ATLAS in the search for  $W' \rightarrow tb$  and improvements to the analysis we are currently working on.
  - Most recent result [ATLAS-CONF-2013-050]
- Why look for  $W' \rightarrow tb$ ?
  - $W'$  is a common feature of many new physics models.
  - The  $W' \rightarrow tb$  channel is complementary to the  $W' \rightarrow \ell\nu$  search channel.
    - When  $W' \rightarrow \ell\nu$  is suppressed, such as when  $m_{\text{neutrino}} > m_{W'}$ , the  $W' \rightarrow tb$  channel has sensitivity competitive with  $W' \rightarrow \ell\nu$ .
    - It is also possible that the  $W'$  is coupled more strongly to the 3<sup>rd</sup> generation quarks. [arXiv:hep-ph/9603349]  
[arXiv:hep-ph/9602390]

# The ATLAS Detector



- Information from all of the ATLAS detector systems is used.
- In the event selection the jets and leptons are required to be centrally located where the detector has the best resolution.

# Event Selection

- Two control regions are formed:
  - Events that fail the  $M(W')$  cut are used to normalize the  $W$ +jets background.
  - Single tagged events are used to verify the modeling of kinematic variables.

Exactly 1 lepton with  $p_T > 30$  GeV.

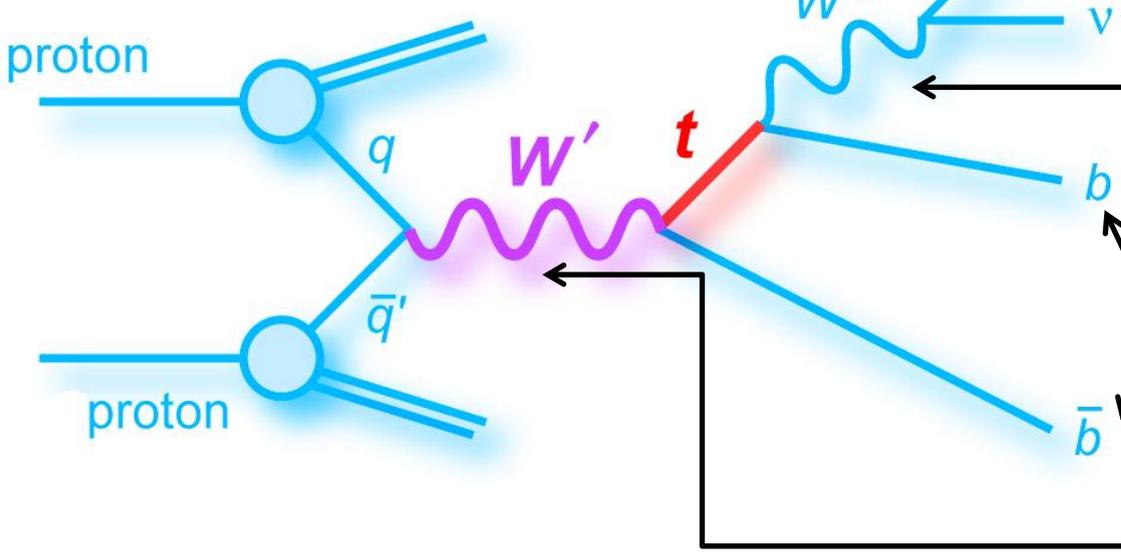
- Electrons:  $|\eta| < 2.47$  excluding  $1.37 < |\eta| < 1.52$
- Muons:  $|\eta| < 2.5$

$MET > 35$  GeV.

Triangular cut  
 $M_T(W) + MET > 60$  GeV.

2 or 3 jets with  $p_T > 25$  GeV,  
 $|\eta| < 2.5$ .  
 Exactly 2 b-tagged jets with  
 70% tagging efficiency.

$M(W') > 270$  GeV.



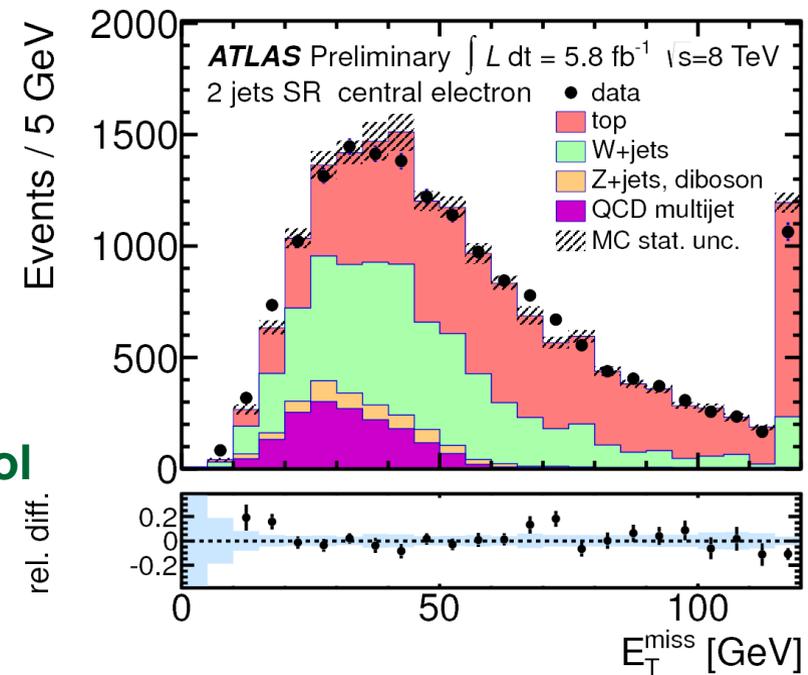
## • Multijet

- Normalized by fitting the MET distribution without the 35 GeV selection cut.
- Shape is obtained by allowing jets to fake electrons in a simulated dijet sample.

## • W+jets

- Normalized by subtracting all other backgrounds from the data in the control region  $M(W') < 270$  GeV for each selection channel.
- Shape is taken from Monte Carlo simulation.

[ATLAS-CONF-2012-132]



- The remainder of the backgrounds are modeled in Monte Carlo simulation and are scaled to the theoretical prediction.

| Sample               | Generator   |
|----------------------|---|
| ttbar                | MC@NLO+HERWIG (Parton Shower and Hadronisation)<br>MC@NLO+JIMMY (Multiple Parton Scattering and Underlying Event) |
| Single-top t-channel | ACERMC+PYTHIA   |
| Single-top Wt        | MC@NLO+HERWIG   |
| Single-top s-channel | MC@NLO+HERWIG   |
| Diboson              | HERWIG  |
| Z+jets               | ALPGEN+HERWIG   |
| W+jets               | ALPGEN+HERWIG   |
| Multijet             | PYTHIA  |
| W'                   | MADGRAPH+PYTHIA   |

- We consider the following uncertainties:
  - **Lepton Energy Scale and Resolution**
    - The uncertainty in the lepton energy scale and energy resolution are propagated through the analysis.
  - **Lepton Identification and Trigger Efficiency**
    - The uncertainty in the lepton identification efficiency and the trigger efficiencies are assessed.
  - **Jet Energy Scale** — One of the dominant systematics
    - The uncertainty in the jet energy scale (JES) is assessed for each jet dependent on its  $p_T$  and position in the detector and is propagated to the calculation of the missing transverse energy.
  - **Jet Energy Resolution**
    - The impact of the jet energy resolution uncertainty is assessed by smearing the jet energy in all samples.
  - **Jet Reconstruction Efficiency**
    - The uncertainty in the jet reconstruction efficiency is assessed by randomly dropping jets from events.

- **b-tagging Performance** — One of the dominant systematics
  - The uncertainty in the b-tagging efficiency and the mistagging rates are estimated from data.
- **Monte Carlo Generator**
  - The uncertainty in the ttbar yield due to the Monte Carlo generator is assessed by taking the larger of the differences between the nominal MC@NLO sample and samples produced using POWHEG+HERWIG and ALPGEN+HERWIG.
- **Parton Shower Modeling**
  - The uncertainty in the ttbar acceptance due to the parton shower modeling is assessed by taking the difference between samples produced using POWHEG+HERWIG and POWHEG+PYTHIA.
- **Initial State Radiation/Final State Radiation (ISR/FSR)**
  - The dependence of the top-quark backgrounds on the ISR/FSR modeling is assessed by taking the largest difference from a set of samples generated using ACERMC+PYTHIA with varying PYTHIA ISR and FSR parameter settings.

## – Parton Distribution Function (PDF)

- The PDF uncertainty is estimated for the top-quark and signal samples normalizing the envelope of CT10, MWST2008NLO68CL, and NNPDF2X at 68% CL to the nominal cross-section.

**Theoretical Cross-section Normalization** – One of the dominant systematics

- For each sample which is normalized to the theoretical cross-section we assess a flat rate uncertainty to account for the scale variations and the uncertainties in  $\alpha_s$  and PDFs.

## – Multijet Background Normalization

- We assess a 50% uncertainty on the multijet background rate.

## – W+jets Background

- By propagating the effects of all other systematics in the control region we assess the uncertainty in the W+jets rate. The uncertainty in the modeling is estimated using samples produced while varying the ALPGEN parameters parton  $p_T$  and the functional form of the factorization scale.

- **Jet Vertex Fraction**
  - The effects of the uncertainty on the jet vertex fraction is assessed on the rates of all samples.
- **Luminosity**
  - The uncertainty in the integrated luminosity used in this analysis is 3.6%, derived from beam-separation scans performed in April 2012.
- **MC Statistics**
  - The impact of limited size simulated samples is taken into account.

- For each channel a **Boosted Decision Tree (BDT)** is trained using the **ROOT Toolkit for Multivariate Data Analysis (TMVA)**.
  - For the training signal the simulation of right handed  $W'$  with a mass of 1750 GeV is used.
  - Variables are selected from a long list of kinematic variables.
  - Selected variables are required to have a discriminating power greater than 20%.
  - Selected variables must be well modeled in the 1-tag control region.

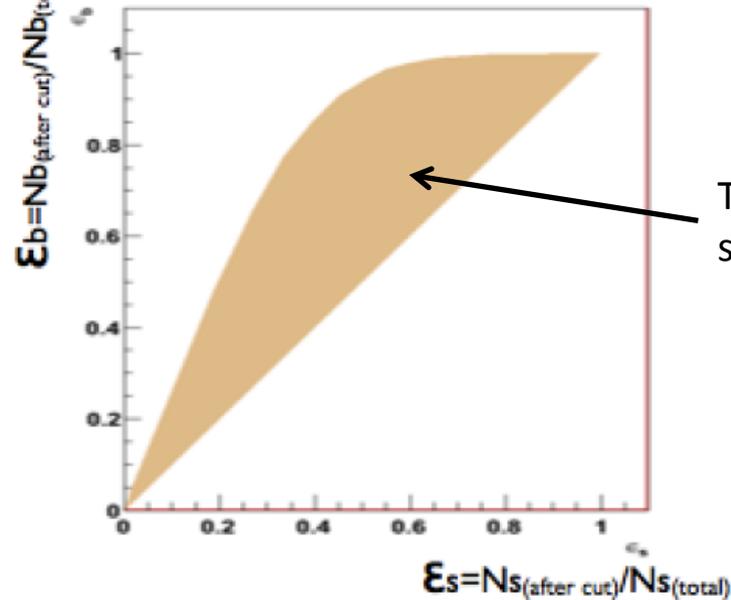
### 2-Jet 2-Tagged BDT Variables

|                                   |                      |
|-----------------------------------|----------------------|
| $m_{tb}$                          | $p_T(b_2)$           |
| $p_T(t)$                          | $p_T(l)$             |
| $\Delta R(l, b_2)$                | $\Delta R(l, t)$     |
| $H_T(l, \text{jets}, \text{MET})$ | $\Delta R(l, W)$     |
| $m_{b_1 b_2}$                     | $\Delta R(b_1, b_2)$ |
| $E_T(W)$                          | sphericity           |
| $p_T(b_1)$                        | aplanarity           |

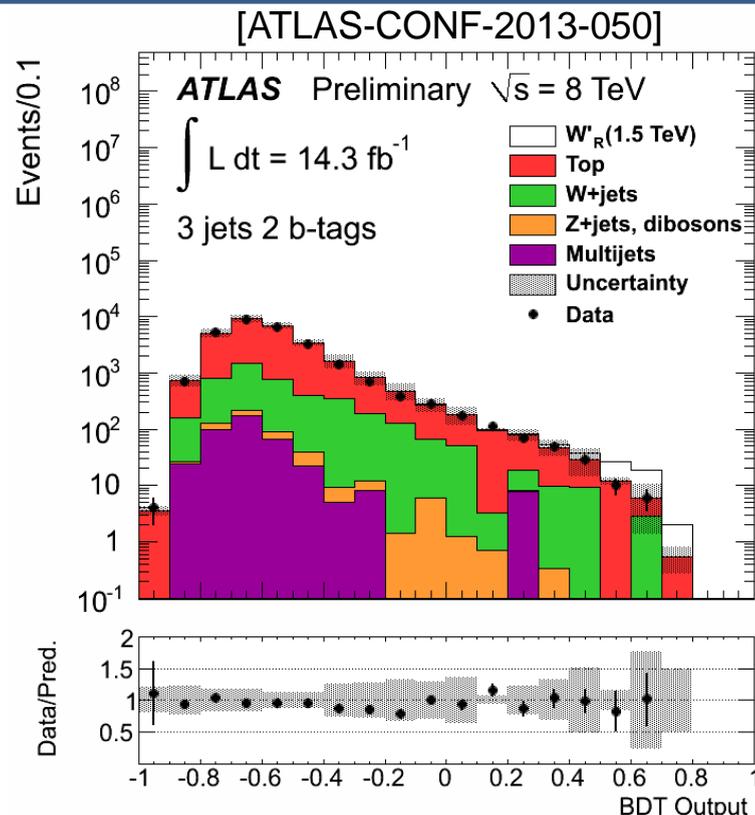
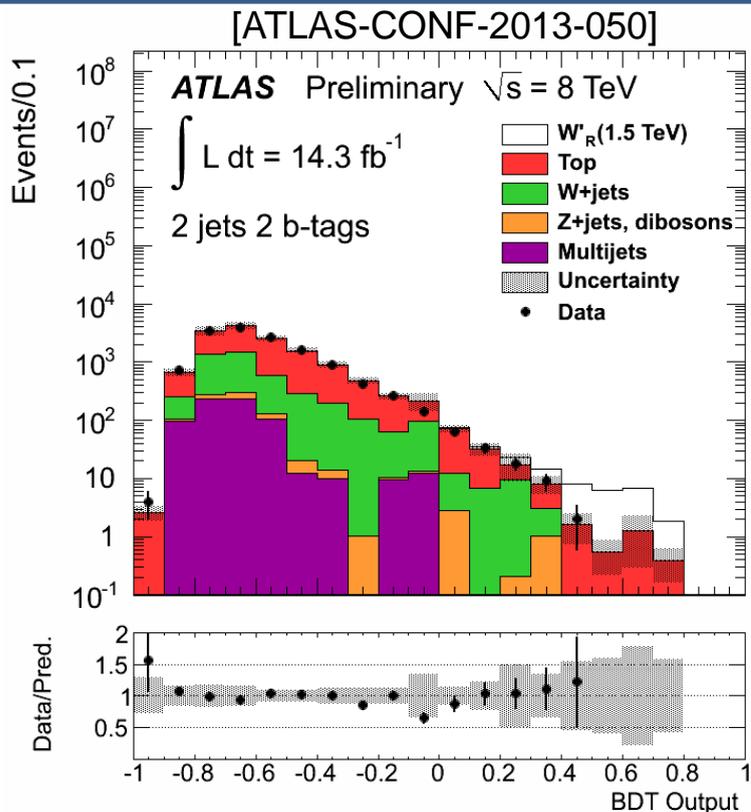
### 3-Jet 2-Tagged BDT Variables

|                    |                    |
|--------------------|--------------------|
| $m_{tb}$           | $\Delta\eta(l, t)$ |
| $p_T(t)$           | $\Delta R(b_1, t)$ |
| sphericity         | $\Delta R(b_1, W)$ |
| $p_T(b_1)$         | $\Delta R(l, b_1)$ |
| $p_T(l)$           | aplanarity         |
| $M_{lb_1 b_2}$     | $\Delta\phi(W, t)$ |
| $\Delta\eta(l, W)$ | -                  |

## Variable Discriminating Power



Discriminating power is determined by plotting the background vs. signal efficiencies ( $\epsilon_B$  vs.  $\epsilon_S$ ) for successive cuts on each variable and taking the area between the resulting curve and the line  $\epsilon_B = \epsilon_S$ .



- The BDT output for the 2 jet and 3 jet selections are above.
- A right handed  $W'$  signal with a mass of 1.5 TeV is added to the backgrounds to illustrate what a signal would look like in the analysis.

# Event Yields

|                         | 2-jet 2-tag channel | 3-jet 2-tag channel |
|-------------------------|---------------------|---------------------|
| $W'_R$ (0.5 TeV)        | $11800 \pm 2700$    | $8200 \pm 1800$     |
| $W'_R$ (1.0 TeV)        | $600 \pm 150$       | $660 \pm 160$       |
| $W'_R$ (1.5 TeV)        | $42 \pm 11$         | $56 \pm 13$         |
| $W'_R$ (2.0 TeV)        | $4.2 \pm 1.1$       | $6.2 \pm 1.5$       |
| $W'_R$ (2.5 TeV)        | $0.69 \pm 0.17$     | $0.87 \pm 0.20$     |
| $W'_R$ (3.0 TeV)        | $0.22 \pm 0.06$     | $0.25 \pm 0.06$     |
| $t\bar{t}$              | $8300 \pm 2100$     | $22000 \pm 5000$    |
| Single-top $t$ -channel | $1000 \pm 270$      | $1400 \pm 400$      |
| Single-top $Wt$         | $400 \pm 80$        | $880 \pm 170$       |
| Single-top $s$ -channel | $310 \pm 90$        | $160 \pm 50$        |
| $W$ +jets               | $3600 \pm 1900$     | $4000 \pm 5000$     |
| Diboson                 | $130 \pm 60$        | $80 \pm 40$         |
| $Z$ +jets               | $26 \pm 20$         | $42 \pm 30$         |
| Multijets               | $710 \pm 350$       | $410 \pm 210$       |
| Total bkg.              | $14400 \pm 3100$    | $29000 \pm 7000$    |
| Data                    | 14138               | 27759               |

- The final event yields for several right handed  $W'$  signal mass points as well as the backgrounds and data.

- With no observed excess in the data we calculate exclusion limits on the production cross-section of the signal as a function of its mass.
  - We use the  $CL_s$  procedure to calculate exclusion limits at the 95% confidence level using a log-likelihood ratio (LLR) as the test statistic.

$$LLR = -2 \ln \frac{\mathcal{L}(data | H_1)}{\mathcal{L}(data | H_0)}$$

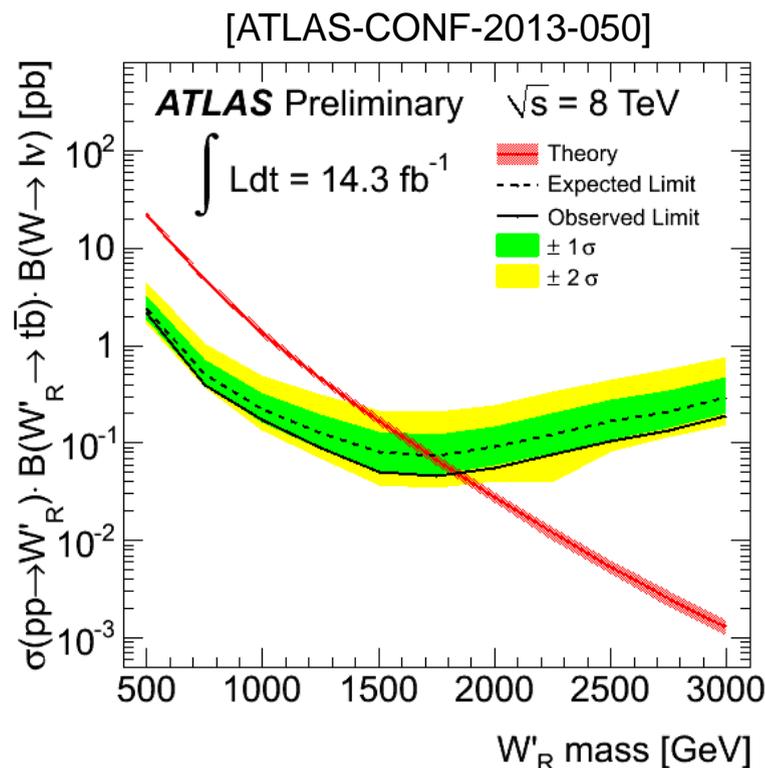
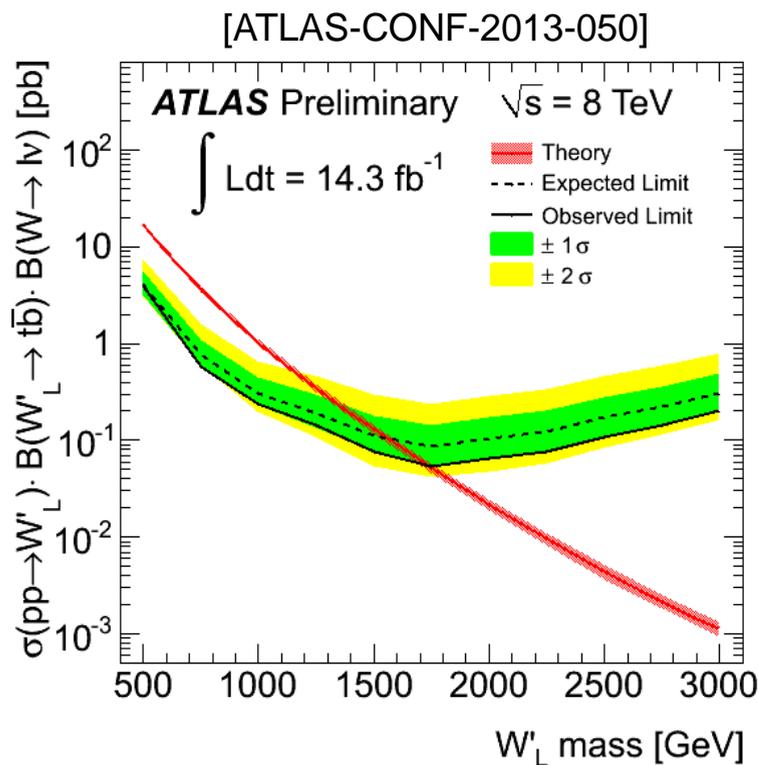
$H_1$  is the test hypothesis which includes the signal

$H_0$  is the null hypothesis which includes no signal

- Pseudo-experiments are generated for each hypothesis.
  - Statistical fluctuations are treated on a bin-by-bin basis as Poisson fluctuations.
  - Systematic uncertainties are treated on a bin-by-bin basis as independent Gaussian fluctuations.

- The bin probabilities are combined into a total probability.
  - Each channel's probability is the product its bins' Poisson probabilities.
  - The total probability is the product of the channels' probabilities.
- $CL_{s+b}$  ( $CL_b$ ) is defined as the fraction of pseudo-experiments generated with the signal-plus-background (background only) hypothesis with LLR greater than the observed or expected LLR.
- Cross-sections are considered excluded at the 95% confidence level if  $CL_s = CL_{s+b}/CL_b < 0.05$ .

- To reduce the effects of the systematics on the limit we fit the  $t\bar{t}b\bar{b}$  yield to data during the statistical analysis.

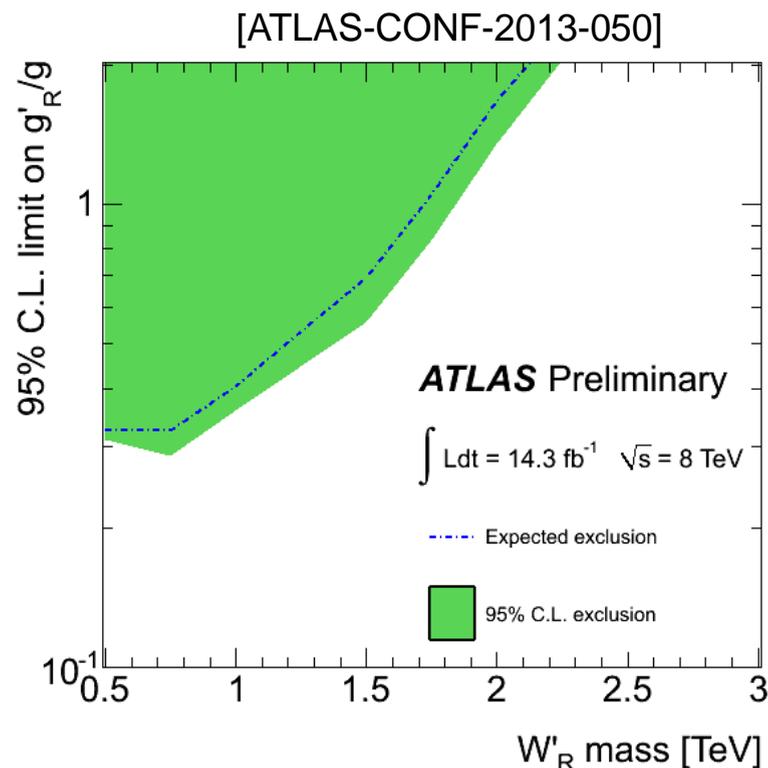
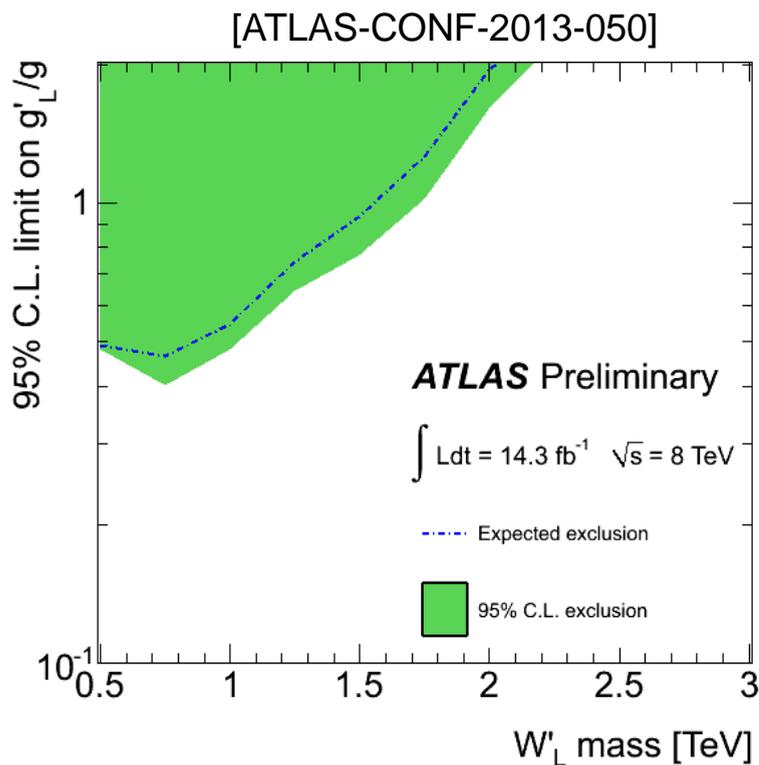


- Both 2 and 3 jet channels are combined and exclusion limits on the cross-section time branching ratio for each  $W'$  mass are placed.
- For a  $W'$  with standard model couplings, the mass lower limits are:
  - $W'_L > 1.74 \text{ TeV}$
  - $W'_R > 1.84 \text{ TeV}$

- Limits can also be set on  $g'$  as a function of  $W'$  mass for models with Lagrangians of the form:

$$\mathcal{L} = \frac{V'_{ij}}{2\sqrt{2}} \bar{f}_i \gamma_\mu \left( g'_{R_{i,j}} (1 + \gamma^5) + g'_{L_{i,j}} (1 - \gamma^5) \right) W'^\mu f_j + h.c.$$

- The production vertex has a  $g'^2$  dependence.
  - By taking the ratio of our cross-section limit and the standard model coupling cross-section for each  $W'$  mass, limits on  $g'/g$  can be derived.
  - The  $W'$  resonance width,  $\Gamma_{W'}$ , is also dependent on  $g'^2$  so the cross section is sensitive to the initial-state quark PDFs.
  - The cross-section's dependence on  $g'/g$  and  $m_{W'}$  is estimated using MADGRAPH and the effects are found to be at most a few percent for  $g'/g < 2$  and thus they are neglected.



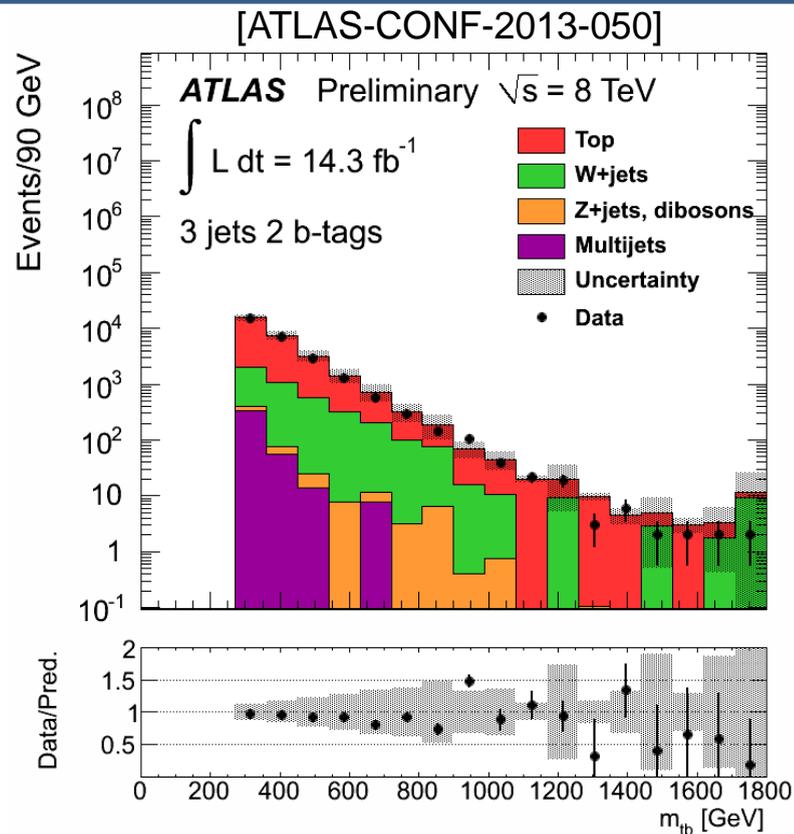
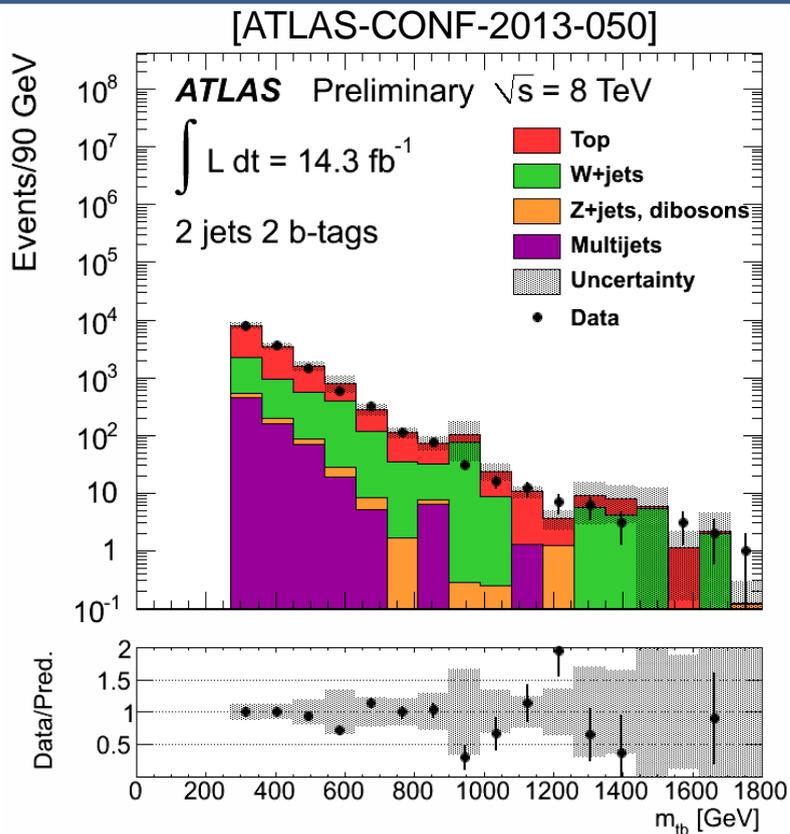
- Limits on  $g'/g$  as functions of  $W'$  mass.
- For values of  $g'/g > 2$ , the  $W'$  width becomes significant and interactions with the initial state quarks' PDF must be taken into account.

- Update with the full  $20.4 \text{ fb}^{-1}$  of 8 TeV data.
  - This includes updated recommendations from other groups.
- Include 1-tag events as an independent channel.
- More robust BDT strategy
  - Training separate BDTs for left and right handed samples, possibly training different BDTs for different mass ranges.
- Improved event selection
  - Relaxing the muon isolation requirement increases the signal efficiency by up to ~50%.
  - Investigating boosted-b-tagging using soft-lepton-tagging. [arXiv:1307.1820]
  - More stringent event selection for 1-tag events may be necessary to control the backgrounds.

- We searched for  $W' \rightarrow tb \rightarrow lvbb$  in  $14.3 \text{ fb}^{-1}$  by considering events with a single lepton, MET, and 2 b-tagged jets and applying boosted decision trees to construct multivariate discriminants.
- When no excess was observed, limits are placed on the mass of a  $W'$  with standard model couplings of:
  - $W'_L > 1.74 \text{ TeV}$
  - $W'_R > 1.84 \text{ TeV}$
- Limits are also calculated for  $g'/g$  as a function of  $W'$  mass.
- Our group also has a hadronic analysis in progress.
- Updated results should be available this fall!

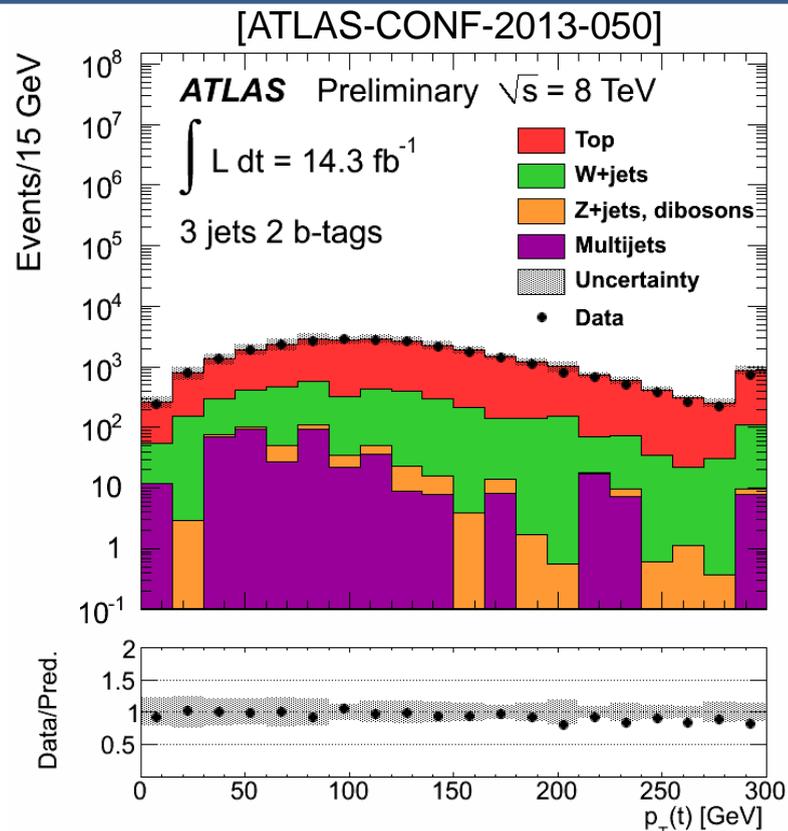
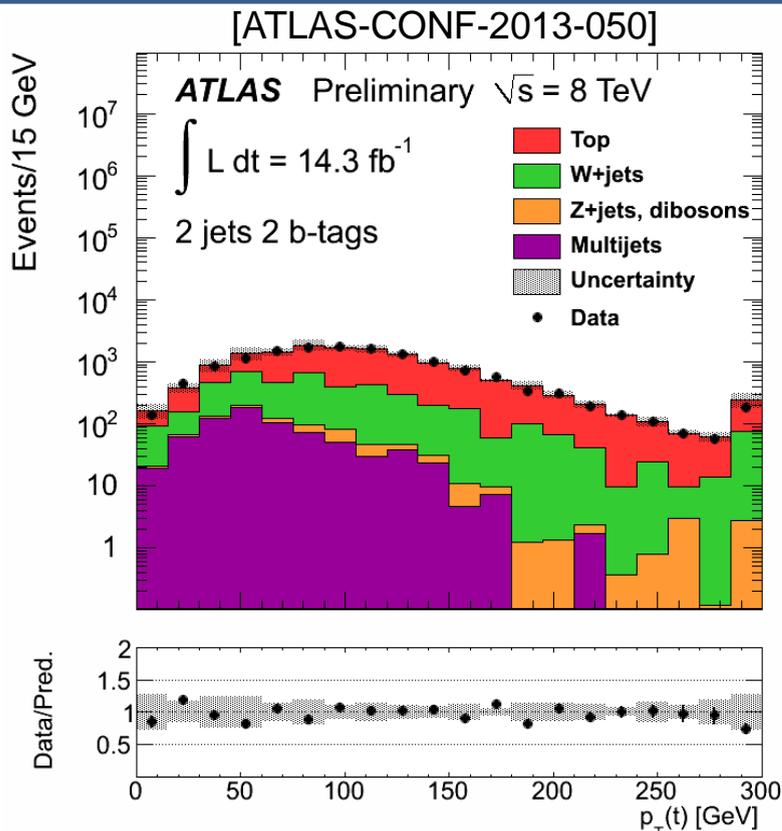
# Backup Slides

# $m_{tb}$ Distribution



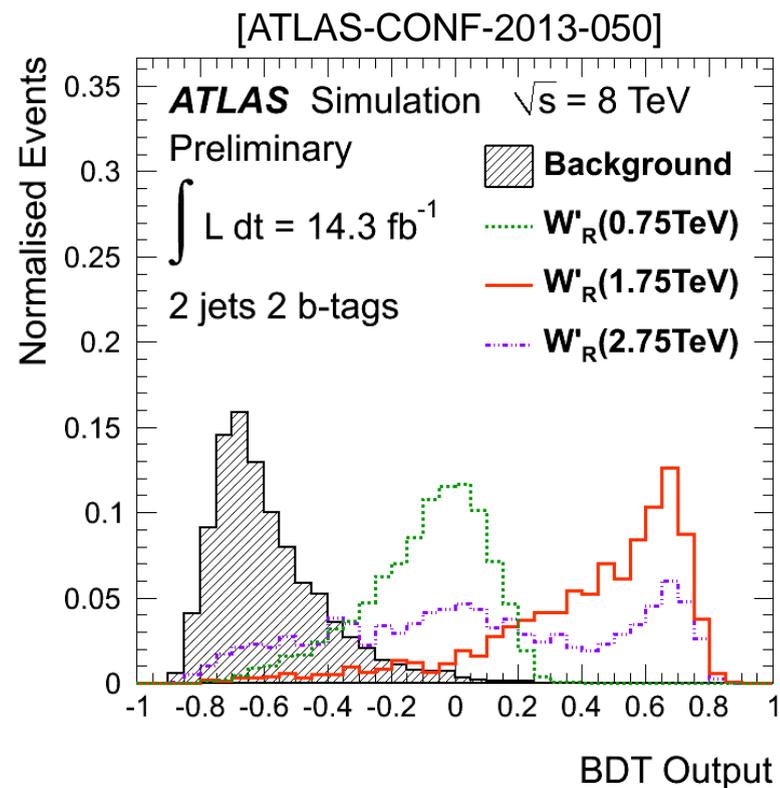
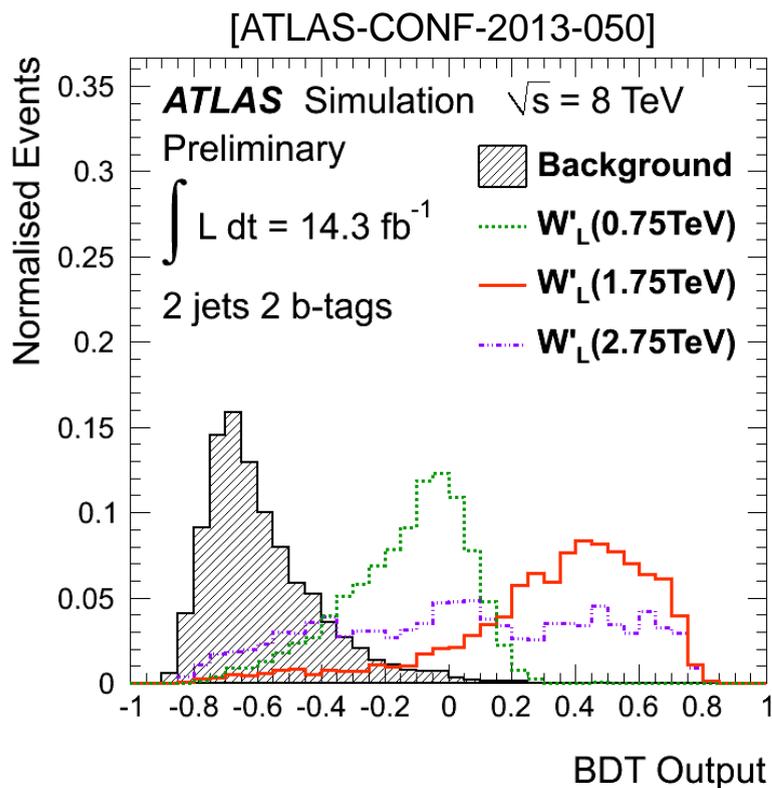
- $m_{tb}$  is the invariant mass of the reconstructed top and bottom quarks, which is the  $W'$  invariant mass.
- This is the most discriminating variable for both selections.

# $p_T(t)$ Distribution



- $p_T(t)$  is the transverse momentum of the reconstructed top quark.
- This is the second most discriminating variable for both selections.

# BDT Distribution 2 Jets



# BDT Distribution 3 Jets

